

Crystal Structure and Properties of $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$

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The crystal structure of $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ was determined using single-crystal and powder X-ray diffraction methods. This phase crystallizes in the monoclinic space group $C2/m$ (No. 12) ($a = 23.324(1)$ Å, $b = 11.388(1)$ Å, $c = 9.8499(3)$ Å, $\beta = 90.104(3)^\circ$; $Z = 2$; $\rho_{\text{calcd.}} = 4.98 \text{ g/cm}^3$), and exhibits a 10-layer structure built from close-packed $[\text{O}(\text{Ba},\text{O})]$ layers with a stacking sequence $(cchhc)_2$. Octahedral sites are occupied by a mixture of Fe^{3+} and Ti^{4+} , with some preferential ordering suggested by analysis of bond valence sums. The structure features vertex-, edge-, and face-sharing of the $[\text{Ti}(\text{Fe})\text{O}_6]$ octahedra. Indexed X-ray powder diffraction data for a polycrystalline specimen are given. $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ and the 8-layer phase $\text{Ba}_4\text{Fe}_2\text{Ti}_{10}\text{O}_{27}$ are built from the same types of polyhedral layers, some of which feature vacant sites between two Ba ions, which substitute for three oxygens in a row. The single-crystal results suggest that the basic structural formula of the phase is $\text{A}_{11}\text{B}_{28}\text{O}_{66+x}$, with the value of x (and hence the Fe/Ti ratio) determined by partial occupancy of one of these vacant sites. Variation of this occupancy factor with synthesis temperature may account for apparent slight differences in the stoichiometry of this phase in polycrystalline and single-crystal form. However, solid solution forma-

tion was not observed for polycrystalline specimens. A comparison of the crystal structure obtained for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ with that previously proposed for “ Zr^{4+} -stabilized $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” indicates that the phase “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” is actually a ternary compound which forms upon addition (either deliberately or inadvertently) of a trivalent ion such as Fe^{3+} or Al^{3+} . The specimens $\text{Ba}_{11}\text{Al}_2\text{Ti}_{26}\text{O}_{66}$, $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Sn}_2\text{O}_{66}$, and $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Zr}_2\text{O}_{66}$ were also prepared and were found to form the $\text{A}_{11}\text{B}_{28}\text{O}_{66+x}$ -type phase. $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ exhibits paramagnetic behavior that deviates somewhat from the Curie–Weiss Law below 75 K. Application of this formalism to the $1/\chi$ vs. T data above 75 K yields an effective moment consistent with the presence of high-spin Fe^{3+} ($S = 5/2$), and a negative Weiss constant (about -25 K) indicating weak cooperative magnetic interactions that are overall antiferromagnetic. The relative permittivity and dielectric loss tangent of a sintered polycrystalline disk were measured at 5.33 GHz, yielding values (corrected for theoretical density) of 55 and $7.7(\pm 0.3) \times 10^{-4}$, respectively.

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Introduction

Ceramic magnetic oxides serve important functions in a wide variety of electronic applications, including wireless communications. For example, circulators and isolators contain ceramic magnets that provide a static magnetic field to interact with the propagating microwave, thus steering and conditioning the signal. The properties required for such ceramics include large saturation magnetization, low dielectric loss, and a high dielectric constant. Materials currently in use include various garnets, spinels, and hexaferrites,^[1,2] which exhibit relatively low relative permittivities (about 15). New materials with larger permittivities would permit miniaturization of these components. To this end, the $\text{BaO}:\text{Fe}_2\text{O}_3:\text{TiO}_2$ phase diagram has been systematically

investigated^[3] to elucidate compounds forming between the high-dielectric-constant polytitanates (e.g. $\text{Ba}_2\text{Ti}_9\text{O}_{19}$ and BaTi_4O_9), and strongly magnetic compounds such as barium hexaferrite ($\text{BaFe}_{12}\text{O}_{19}$). The phase diagram study confirmed the formation of a number of ternary phases with new structural types, including the title compound, denoted previously as phase “G”^[3] with a stoichiometry close to “ $\text{Ba}_{14}\text{Fe}_2\text{Ti}_{35}\text{O}_{87}$ ” ($= 0.28:0.02:0.70 \text{ BaO}:\text{Fe}_2\text{O}_3:\text{TiO}_2$), as estimated by the disappearing phase method. The present report describes the crystal structure and more precise stoichiometry of this compound, as determined by single-crystal and X-ray powder diffraction methods, and characterization of its dielectric and magnetic properties. In addition, the relationship between this ternary phase and the compound reported as “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” is described.

Results and Discussion

Description of the Crystal Structure and Stoichiometry

The results of the single-crystal structure determination are given in Table 1. Selected bond lengths and bond-val-

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ence sums are collected in Table 2, and indexed experimental X-ray powder diffraction data for a polycrystalline specimen are given in Table 3. As seen in Table 1, electron density was observed at the O22 site indicating a partial occupancy factor of 0.24(7), resulting in an overall stoichiometry of Ba₁₁FeTi₂₇O_{66.5} (BaO:Fe₂O₃:TiO₂ = 0.2857:0.0130:0.7013) (*Z* = 2). With site O22 empty, the overall stoichiometry would be Ba₁₁Fe₂Ti₂₆O₆₆ (= 0.2895:0.0263:0.6842 BaO:Fe₂O₃:TiO₂). The basic structural formula for this phase is therefore A₁₁B₂₈O_{66+x}, where A = Ba²⁺ and B = Ti⁴⁺, substituted with a small amount (<10 atom%) of a trivalent metal which determines the value of *x*. The subsolidus study^[3] of the system at 1260 °C suggested little, if any, solid solution formation for this phase, and several closely spaced specimens indicated that the phase forms near the stoichiometry 0.28:0.02:0.70 BaO:Fe₂O₃:TiO₂, which is closer to the first formula with less Fe and a higher oxygen content. However, two polycrystalline specimens subsequently prepared precisely at Ba₁₁FeTi₂₇O_{66.5} and Ba₁₁Fe₂Ti₂₆O₆₆ were not single-phase, despite repeatedly heating for hundreds of hours. The former composition contained small amounts (several %) of BaTi₂O₅ and Ba₆Ti₁₇O₄₀, while the latter contained an appreciable amount (>10%) of Ba₄Fe₂Ti₁₀O₂₇. According to the subsolidus phase diagram,^[3] this suggests that at 1260° the compound forms with a composition between Ba₁₁FeTi₂₇O_{66.5} and Ba₁₁Fe₂Ti₂₆O₆₆, with a lower O22 site occupancy factor than that determined for crystals grown by cooling from 1390 °C. As discussed below, these slight variations in Fe/Ti/O stoichiometry (note, however, that the A-cation/B-cation ratios remain constant) are readily accommodated by the crystal structure, and have been observed for other compounds occurring in the BaO:Fe₂O₃:TiO₂ system.^[3–5] For convenience, the compound in the present study is referred to as Ba₁₁FeTi₂₇O_{66.5}, according to Table 1. The X-ray powder pattern, calculated using the single-crystal results (Table 1) and the refined unit cell (Table 3), is in excellent agreement with the pattern observed for a polycrystalline sample (0.28:0.02:0.70 BaO:Fe₂O₃:TiO₂) of this phase, as shown in Figure 1.

The structure of Ba₁₁FeTi₂₇O_{66.5} may be described as a 10-layer (10*L*) close-packed (*cp*) arrangement built from [O] and [Ba,O] layers, stacked along the *a*-axis, with the transition metals occupying octahedral interstices. The connectivity of the structure includes vertex-, edge-, and face-sharing of [(Ti⁴⁺/Fe³⁺)O₆] octahedra. The *a*-parameter of the unit cell indicates an average *cp* layer thickness of 2.362 Å, which is similar to that observed for other close-packed structures in this chemical system.^[6,7] The arrangement is illustrated in Figure 2, as viewed along the ⟨011⟩ direction to reveal the stacking pattern of the cubic (*ccp*), and hexagonal (*hcp*) close-packed layers. The stacking sequence is (*cchhc*)₂ (*c* denotes a cubic close-packed layer and *h* denotes a hexagonally close-packed layer), where octahedral interstices share faces across *hcp* layers. In pairs of face-sharing octahedra that are occupied, the cations are displaced away from each other resulting in cation-cation distances of the order of 2.7 Å to 2.8 Å.

Table 1. Structural Parameters obtained for Ba₁₁FeTi₂₇O_{66.5}, M = Ti⁴⁺(Fe³⁺)

Atom	Site	x	y	z	100*U _{iso}	occ.
Ba 1	2a	0	0	0	1.58(12)	
Ba 2	4h	1/2	0.22715(23)	1/2	1.61(8)	
Ba 3	4i	0.89468(9)	0	0.67147(22)	1.01(7)	
Ba 4	8j	0.60238(7)	0.22859(16)	0.83942(16)	1.51(6)	
Ba 5	4i	0.78713(10)	0	0.33795(24)	1.46(8)	
M 1	8j	0.34924(21)	0.1378(5)	0.7628(5)	1.48(15)	
M 2	8j	0.45461(18)	0.2624(4)	0.8370(4)	1.03(13)	
M 3	4i	0.0494(3)	0	0.6582(7)	0.94(20)	
M 4	4i	0.3580(3)	0	0.4872(7)	1.36(23)	
M 5	8j	0.74963(19)	0.1340(5)	0.7190(5)	1.37(15)	
M 6	4i	0.1401(3)	0	0.9918(7)	1.01(21)	
M 7	4i	0.2617(3)	0	0.0032(7)	1.10(21)	
M 8	8j	0.14852(20)	0.7558(4)	0.5063(5)	1.13(15)	
M 9	4f	1/4	3/4	0	1.50(24)	
M10	4i	0.5373(3)	0	0.1630(8)	2.1(3)	
O1	4i	0.6040(10)	0	0.3246(23)	0.66(6)	
O2	8j	0.5024(7)	0.1329(15)	0.2383(16)	0.66(6)	
O3	4g	0	0.2386(21)	0	0.66(6)	
O4	2c	0	0	1/2	0.66(6)	
O5	4i	0.8988(10)	0	0.1725(23)	0.66(6)	
O6	8j	−0.0025(7)	0.1224(14)	0.2513(16)	0.66(6)	
O7	8j	0.3988(7)	0.1292(15)	0.4214(16)	0.66(6)	
O8	8j	0.1013(7)	0.2441(15)	0.3195(16)	0.66(6)	
O9	8j	0.9019(7)	0.1310(15)	0.4211(16)	0.66(6)	
O10	8j	0.1986(7)	0.1326(15)	0.4192(16)	0.66(6)	
O11	8j	0.2100(7)	0.2511(15)	0.1709(16)	0.66(6)	
O12	8j	0.5962(7)	0.1267(14)	0.0957(16)	0.66(6)	
O13	8j	0.2979(7)	0.1300(15)	0.0900(16)	0.66(6)	
O14	8j	0.0970(7)	0.1203(15)	0.0786(16)	0.66(6)	
O15	8j	0.6889(7)	0.1256(14)	0.4118(16)	0.66(6)	
O16	4i	0.2937(10)	0	0.3369(23)	0.66(6)	
O17	4i	0.6973(10)	0	0.1673(23)	0.66(6)	
O18	8j	0.3044(7)	0.2488(15)	0.3389(16)	0.66(6)	
O19	4i	0.2099(10)	0	0.1806(23)	0.66(6)	
O20	8j	0.7998(7)	0.1130 (15)	0.0609(16)	0.66(6)	
O21	2b	1/2	0	0	0.66(6)	
O22	4i	0.401(4)	0	0.210(10)	0.66(6)	0.24(7)

The polyhedral arrangements within each of the ten layers in the Ba₁₁FeTi₂₇O_{66.5} structure are shown in Figure 3. This compound and Ba₄Fe₂Ti₁₀O₂₇^[6] (isostructural with Ba₄Al₂Ti₁₀O₂₇^[8]) are built from the same three types of layers: the first type contains strings of three edge-sharing octahedra plus isolated octahedra (layers 1,5,6,10), the second contains groups of five edge-sharing octahedra plus isolated octahedra (layers 2,4,7,9), and a third type contains continuous zigzag strips of edge-sharing octahedra (layers 3,8). The 8*L* phase Ba₄Fe₂Ti₁₀O₂₇^[6] is built from two layers of the first type, four layers of the second type, and two layers of the third type (Figure 3 of ref.^[6]). Not surprisingly, these two phases occur in equilibrium with one another.^[3] The similarity of the two structures suggests that intergrowths between them may form easily, and could be the source of the relatively high residual in the structure refinement.

An interesting crystal-chemical feature, reminiscent of the barium polytitanates,^[9] is the occurrence of vacancies in the oxygen/barium *cp* layers. As seen in Figure 3, and shown in detail in Figure 4, the vacancies occur in the first

Table 2. Selected bond lengths and cation bond valence sums (BVS)* for Ba₁₁FeTi₂₇O_{66.5}; * calculated using the formalism and parameters ($B = 0.37$; r_0 values 2.285, 1.815, 1.759 for Ba²⁺, Ti⁴⁺, Fe³⁺, respectively) from ref.^[23]

Ba 1	–O3	2.718 (24) Å	(2 ×)	M 3	–O4	1.939 (16) Å	
	–O5	2.912 (23) Å	(2 ×)		–O5	2.060 (24) Å	
	–O6	2.844 (16) Å	(4 ×)		–O6	1.985 (17) Å	(2 ×)
	–O14	2.757 (16) Å	(4 ×)		–O9	2.032 (17) Å	(2 ×)
	BVS	3.00			BVS(Ti ⁴⁺)	3.62	
Ba 2				M 4	BVS(Fe ³⁺)	3.11	
	–O2	2.795 (16) Å	(2 ×)		–O1	2.056 (24) Å	
	–O4	3.108 (16) Å			–O7	1.870 (17) Å	(2 ×)
	–O6	2.992 (16) Å	(2 ×)		–O15	2.059 (17) Å	(2 ×)
	–O7	2.723 (16) Å	(2 ×)		–O16	2.109 (24) Å	
	–O8	2.979 (16) Å	(2 ×)		BVS(Ti ⁴⁺)	3.75	
	–O9	2.908 (16) Å	(2 ×)		BVS(Fe ³⁺)	3.33	
Ba 3	BVS	2.21		M 5			
	–O4	2.985 (16) Å			–O10	1.822 (17) Å	
	–O6	2.976 (16) Å	(2 ×)		–O11	1.959 (17) Å	
	–O8	2.784 (16) Å	(2 ×)		–O13	2.187 (17) Å	
	–O9	2.890 (16) Å	(2 ×)		–O16	1.911 (16) Å	
	–O10	2.796 (16) Å	(2 ×)		–O18	1.904 (17) Å	
	–O14	2.826 (16) Å	(2 ×)		–O19	2.050 (17) Å	
	–O19	2.844 (23) Å			BVS(Ti ⁴⁺)	4.12	
Ba 4	BVS	2.57			BVS(Fe ³⁺)	3.54	
	–O2	2.780 (16) Å		M 6	–O5	1.858 (24) Å	
	–O3	2.892 (16) Å			–O14	1.904 (17) Å	(2 ×)
	–O5	3.094 (16) Å			–O19	2.472 (24) Å	
	–O6	3.019 (16) Å			–O20	1.973 (17) Å	(2 ×)
	–O7	2.810 (16) Å			BVS(Ti ⁴⁺)	3.95	
	–O9	3.027 (16) Å			BVS(Fe ³⁺)	3.40	
	–O12	2.785 (16) Å					
	–O13	2.675 (16) Å			–O13	1.908 (17) Å	(2 ×)
	–O14	2.922 (16) Å			–O17	1.935 (24) Å	
	–O18	2.809 (16) Å			–O19	2.128 (24) Å	
	–O20	3.070 (16) Å			–O20	2.029 (17) Å	(2 ×)
	–O22	2.650 (16) Å			BVS(Ti ⁴⁺)	3.85	
	BVS	2.61			BVS(Fe ³⁺)	3.31	
Ba 5	–O5	3.076 (22) Å		M 8	–O7	1.953 (17) Å	
	–O9	3.174 (16) Å	(2 ×)		–O8	2.146 (17) Å	
	–O10	2.851 (16) Å	(2 ×)		–O9	1.887 (17) Å	
	–O15	2.799 (16) Å	(2 ×)		–O10	1.928 (17) Å	
	–O17	2.686 (16) Å			–O15	1.990 (17) Å	
	–O18	2.889 (16) Å	(2 ×)		–O18	1.882 (17) Å	
	–O20	3.034 (16) Å	(2 ×)		BVS(Ti ⁴⁺)	4.13	
	BVS	2.23			BVS(Fe ³⁺)	3.55	
M 1				M 9	–O11	1.927 (16) Å	(2 ×)
	–O1	2.098 (16) Å			–O13	1.977 (17) Å	(2 ×)
	–O8	1.949 (16) Å			–O20	2.037 (17) Å	(2 ×)
	–O11	1.986 (16) Å			BVS(Ti ⁴⁺)	3.88	
	–O12	1.893 (16) Å			BVS(Fe ³⁺)	3.34	
	–O15	1.943 (16) Å		M10	–O1	2.226 (24) Å	
	–O17	2.031 (16) Å			–O2	1.879 (17) Å	(2 ×)
M 2	BVS(Ti ⁴⁺)	3.88			–O12	2.100 (17) Å	(2 ×)
	BVS(Fe ³⁺)	3.34			–O21	1.827 (13) Å	
	–O2	1.923 (16) Å			BVS(Ti ⁴⁺)	3.94	
	–O3	1.924 (16) Å			BVS(Fe ³⁺)	3.39	
	–O6	1.932 (16) Å					
	–O8	2.021 (17) Å					
	–O12	2.058 (17) Å					
	–O14	1.982 (17) Å					
	BVS(Ti ⁴⁺)	3.94					
	BVS(Fe ³⁺)	3.39					

Table 3. X-ray powder diffraction data for Ba₁₁FeTi₂₇O_{66.5}; *C2/m*; *a* = 23.324(1) Å, *b* = 11.388(1) Å, *c* = 9.8499(3) Å, β = 90.104(3)°

<i>h</i>	<i>k</i>	<i>l</i>	$2\theta_{\text{obs}}$	<i>I</i> _{obs}	$2\theta_{\text{calcd.}}$	$\Delta 2\theta$	<i>d</i> _{obs.}
1	1	0	8.633	4	8.634	−0.001	10.23
−1	1	1	12.460	<1	12.458	0.002	7.098
0	2	0	15.569	1	15.550	0.019	5.686
2	2	0	17.332	6	17.317	0.015	5.112
−4	0	1	17.637	2	17.647	−0.010	5.024
2	2	1	19.548	17	19.541	0.007	4.537
4	2	0	21.800	1	21.798	0.002	4.0735
5	1	1	22.481	<1	22.483	−0.002	3.9516
3	1	2	22.751	4	22.753	−0.002	3.9053
4	2	1	23.622	25	23.625	−0.003	3.7632
6	0	1	24.603	9	24.615	−0.012	3.6153
2	2	2	25.082	5	25.086	−0.004	3.5474
1	3	1	25.414	3	25.416	−0.002	3.5018
3	3	0	26.120	2	26.102	0.018	3.4087
5	1	2	27.460	<1	27.476	−0.016	3.2453
3	3	1	27.668	5	27.659	0.009	3.2214
7	1	0	27.867	<1	27.875	−0.008	3.1989
2	0	3	28.216	19	28.227	−0.011	3.1601
1	1	3	28.537	27	28.534	0.003	3.1253
6	2	1	29.249	100	29.246	0.003	3.0508
8	0	0	30.640	13	30.639	0.001	2.9154
0	2	3	31.429	84	31.426	0.003	2.8440
−5	3	1	31.688	25	31.682	0.006	2.8213
3	3	2	31.911	36	31.900	0.011	2.8021
−2	2	3	32.369	10	32.360	0.009	2.7635
−7	1	2	33.352	7	33.353	−0.001	2.6842
−2	4	1	33.623	2	33.626	−0.003	2.6632
5	1	3	34.300	2	34.315	−0.015	2.6122
8	2	0	34.526	1	34.534	−0.008	2.5956
−4	2	3	35.051	1	35.051	0.000	2.5579
−5	3	2	35.477	9	35.474	0.003	2.5282
8	2	1	35.753	20	35.766	−0.013	2.5093
−4	4	1	36.244	6	36.241	0.003	2.4764
−1	3	3	36.347	6	36.347	0.000	2.4696
−7	3	1	36.997	3	36.998	−0.001	2.4277
−2	0	4	37.273	7	37.276	−0.003	2.4104
−3	3	3	37.996	2	37.992	0.004	2.3662
10	0	0	38.570	12	38.568	0.002	2.3323
3	1	4	39.167	8	39.168	−0.001	2.2981
4	4	2	39.686	49	39.681	0.005	2.2692
−9	1	2	40.036	3	40.042	−0.006	2.2502
6	4	1	40.293	8	40.295	−0.002	2.2364
−2	2	4	40.618	<1	40.613	0.005	2.2193
−5	3	3	41.130	2	41.122	0.008	2.1928
3	5	0	41.252	1	41.270	−0.018	2.1866
9	3	0	42.191	9	42.187	0.004	2.1401
−3	5	1	42.308	13	42.317	−0.009	2.1344
2	4	3	42.737	8	42.720	0.017	2.1140
−10	2	1	42.840	7	42.843	−0.003	2.1091
4	2	4	42.925	6	42.904	0.021	2.1052
6	4	−2	43.426	87	43.418	0.008	2.0820
−1	5	2	43.926	<1	43.932	−0.006	2.0595
−8	2	3	44.432	6	44.425	0.007	2.0372
8	2	3	44.519	4	44.499	0.020	2.0334
4	4	3	44.892	1	44.896	−0.004	2.0174
5	5	1	45.227	11	45.221	0.006	2.0032
−3	3	4	45.378	14	45.369	0.009	1.9969
−6	2	4	46.387	2	46.400	−0.013	1.9558
7	1	4	46.525	3	46.548	−0.023	1.9503
12	0	0	46.692	1	46.694	−0.002	1.9437
−11	1	2	47.295	<1	47.281	0.014	1.9203
8	4	−2	48.296	47	48.290	0.006	1.8828

Table 3. (Continued)

<i>h</i>	<i>k</i>	<i>l</i>	$2\theta_{\text{obs}}$	<i>I</i> _{obs}	$2\theta_{\text{calcd.}}$	$\Delta 2\theta$	<i>d</i> _{obs.}
8	0	4	48.395	41	48.387	0.008	1.8792
−1	5	3	48.775	15	48.774	0.001	1.8655
7	5	1	49.302	9	49.308	−0.006	1.8468
2	6	1	49.490	11	49.490	0.000	1.8402
11	3	1	50.121	1	50.113	0.008	1.8185
12	2	1	50.437	24	50.444	−0.007	1.8078
−10	2	3	50.526	22	50.524	0.002	1.8048
−5	1	5	50.892	3	50.876	0.016	1.7927
4	2	5	51.510	4	51.513	−0.003	1.7727
7	3	4	52.082	6	52.082	0.000	1.7545
2	6	2	52.224	4	52.204	0.020	1.7501
−13	1	1	52.394	6	52.394	0.000	1.7448
11	3	2	52.811	2	52.816	−0.005	1.7320
4	6	2	54.082	18	54.084	−0.002	1.6943
−6	2	5	54.580	35	54.574	0.006	1.6800
1	5	4	55.032	1	55.021	0.011	1.6672
14	0	1	55.947	19	55.946	0.001	1.6421
2	6	3	56.530	2	56.519	0.011	1.6266
1	7	0	56.683	3	56.671	0.012	1.6225
−9	5	2	56.890	3	56.910	−0.020	1.6171
9	3	4	57.011	3	57.004	0.007	1.6140
12	4	0	57.342	1	57.352	−0.010	1.6054
−1	7	1	57.505	1	57.505	0.000	1.6013
3	1	6	57.954	1	57.957	−0.003	1.5899
−14	2	1	58.408	11	58.406	0.002	1.5787
−5	5	4	58.580	9	58.575	0.005	1.5744
−8	6	1	58.735	8	58.742	−0.007	1.5706
−13	1	3	59.132	3	59.145	−0.013	1.5610
−7	3	5	59.620	1	59.626	−0.006	1.5494
−13	3	2	59.961	1	59.962	−0.001	1.5414
−12	4	2	60.597	24	60.593	0.004	1.5268
−4	2	6	60.758	18	60.750	0.008	1.5231
−5	7	1	60.992	12	60.991	0.001	1.5178
−3	7	2	61.106	12	61.107	−0.001	1.5153
6	0	6	61.262	4	61.284	−0.022	1.5118
−1	3	6	61.612	2	61.624	−0.012	1.5040
−7	5	4	62.001	2	62.019	−0.018	1.4955
−11	3	4	62.688	1	62.688	0.000	1.4808
−15	1	2	63.139	<1	63.148	−0.009	1.4713
−5	7	2	63.361	2	63.362	−0.001	1.4666
3	5	5	63.535	3	63.542	−0.007	1.4631
7	1	6	63.717	1	63.714	0.003	1.4593
4	6	4	63.915	1	63.917	−0.002	1.4553
7	7	1	64.391	<1	64.382	0.009	1.4457
13	1	4	64.806	1	64.815	−0.009	1.4374
−8	6	3	65.072	2	65.080	−0.008	1.4322
8	0	6	65.234	6	65.218	0.016	1.4290
0	4	6	65.570	48	65.590	−0.020	1.4225
8	4	5	65.937	6	65.955	−0.018	1.4155
11	5	3	66.647	6	66.634	0.013	1.4021
2	8	1	66.838	9	66.836	0.002	1.3985
−15	1	3	66.986	6	66.994	−0.008	1.3958
−8	2	6	67.420	<1	67.395	0.025	1.3879
15	3	2	67.848	17	67.829	0.019	1.3802
9	7	0	67.966	14	67.980	−0.014	1.3780
−3	1	7	68.145	5	68.144	0.001	1.3749
−4	8	1	68.428	1	68.450	−0.022	1.3699
−7	5	5	68.923	5	68.947	−0.024	1.3612
−2	6	5	69.140	5	69.129	0.011	1.3575
13	3	4	69.386	1	69.370	0.016	1.3533
−12	4	4	69.818	8	69.837	−0.019	1.3460
12	2	5	69.972	11	69.978	−0.006	1.3434

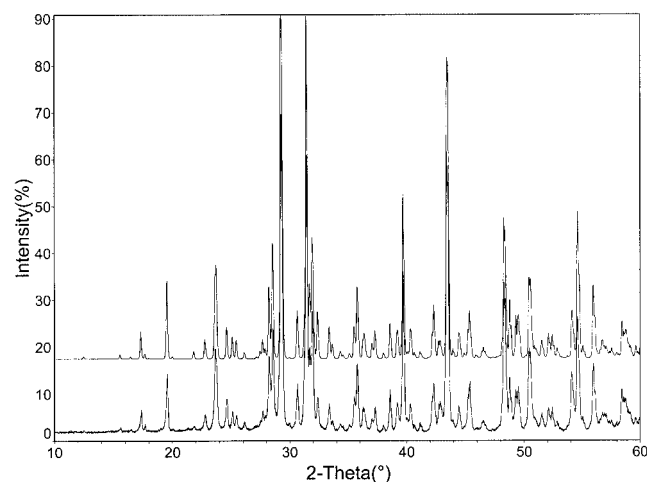


Figure 1. The X-ray powder diffraction pattern for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ calculated from the single-crystal results (*top*) is in excellent agreement with the observed pattern (*bottom*) for a polycrystalline sample ($\text{BaO}:\text{Fe}_2\text{O}_3:\text{TiO}_2 = 0.28:0.02:0.70$); the stoichiometry of this phase in polycrystalline form (prepared at 1260°) apparently differs slightly from that of the single crystals (grown by cooling from 1390°), suggesting slightly different occupancy factors for the oxygen site O22 (Table 2); charge-balance is maintained by a corresponding change in the Fe/Ti ratio; in both cases, however, the A-cation/B-cation ratio remains the same (11:28)

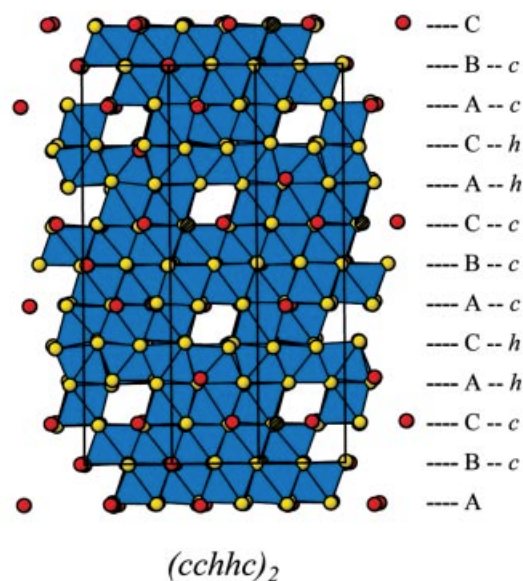


Figure 2. The crystal structure of $10L \text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ ($C2/m$; $a = 23.324(1)$, $b = 11.388(1)$, $c = 9.8499(3)$ Å; $\beta = 90.104(3)^\circ$) viewed approximately parallel to the close-packed (*cp*) $[\text{O},(\text{Ba},\text{O})]$ layers; the 10-layer stacking sequence is $(cchhc)_2$; the octahedra are occupied by $\text{Ti}^{4+}/\text{Fe}^{3+}$, with some preferential ordering suggested by analysis of bond valence sums; red spheres represent Ba^{2+} , yellow spheres represent oxygen with the hatched yellow spheres indicating the partially occupied O22 site (Table 2)

and second types of layers, where two Ba ions replace three oxygens in a row, reducing their coordination number from twelve to eleven. For $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$, this results in a coordination number of 11 and a relatively low bond-valence sum for Ba2 (Table 2, Figure 4). The partially occupied site

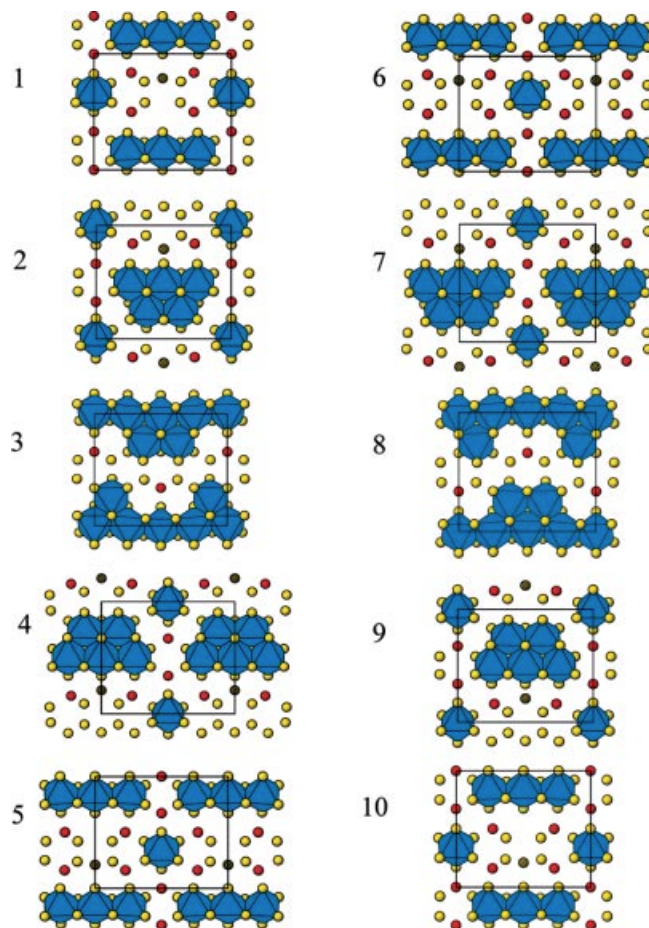


Figure 3. The structural arrangements within each of the ten layers of the close-packed structure of $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$, viewed perpendicular to the layers; the representations are the same as in Figure 1; as seen here, the structure is built from three types of layers (*type 1* = layers 1,5,6,10; *type 2* = layers 2,4,7,9; *type 3* = layers 3,8); the same types of layers occur in the $8L$ structure of $\text{Ba}_4\text{Fe}_2\text{Ti}_{10}\text{O}_{27}$; as observed for the barium polytitanate phases, the *cp* layers (types 1 and 2) are characterized by the occurrence of vacancies, wherein two Ba ions substitute for three oxygen positions in a row; the partially occupied O22 site (hatched yellow spheres) corresponds to one of these vacancies between two Ba ions

O22 corresponds to one of these *cp* vacancies, and occurs within bonding distance of Ba4 (Table 2, Figure 4). The occupancy factor of 0.24(7) for O22 obtained in the single-crystal study therefore suggests that 25% of the Ba4 sites are 12-coordinated and 75% are 11-coordinated.

The bond-valence sums about the M sites (Table 2) were calculated separately using the r_0 parameters for Ti^{4+} and Fe^{3+} , and are generally consistent with the high Ti-content of this phase. However, site M3 was observed to have slightly longer bond lengths than the other sites, and returns a bond-valence sum suggesting preferential occupation by the slightly larger trivalent Fe ion (Table 3). The M3 sites occur in pairs of vertex-linked octahedra in adjacent layers of the first type described above (Figure 3, layers 1,5,6,10,7). Interestingly, the coordination sphere about M3 has no next-nearest neighbors within 3.5 Å, whereas the other M sites have from 3 to 9, with an average of 7, next-

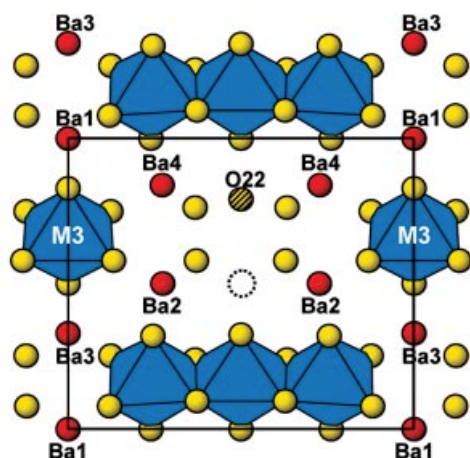


Figure 4. Magnified view of layer 1 in the structure of $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ (representations as in Figure 1), labeled to clarify the positions of the vacancies in the cp layers (dotted circle and hatched yellow O22 site); using the ligands on the upper and lower faces of the octahedra as a guide, the Ba2 sites are seen to occur in the lower cp layer, with the vacancy (dotted sphere) between them, thereby reducing their coordination number to eleven (Table 3); similarly, the Ba4 sites are seen to occur in the upper cp layer, but in this case the vacant site between them appears to be partially occupied (O22, Table 2), according to the electron density observed in the single-crystal study; note the marked displacement of O22 out of the cp row towards the interior of the unit cell; sites M3 occur in vertex-linked pairs of octahedra in adjacent layers, and exhibit bond-valence sums suggestive of preferential occupation by the somewhat larger Fe^{3+} ion

nearest neighbors within this distance. The M3 sites are therefore unusually isolated in the structure, which may facilitate accommodation of a larger cation.

Relationship to “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” (A Ternary Compound)

The synthesis, stability, and crystal chemistry of the phase referred to as “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” has long been of interest to researchers concerned with the phases occurring between BaTiO_3 and TiO_2 , the so-called “barium polytitanates”, which have technically important dielectric properties. The existence of $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ was first reported by Jonker and Kwestroo^[10] during the course of phase equilibria studies of the ternary systems with either ZrO_2 or SnO_2 . They reported considerable solid solution of both oxides in the phase labeled $\text{Ba}_2\text{Ti}_5\text{O}_{12}$. However, many studies of the pure BaO-TiO_2 system failed to confirm the formation of this binary phase.^[11–14] A subsequent study^[15] found that “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” did not form in the ternary system with SnO_2 , in contrast to the earlier report.^[10] The phase did form, however, upon addition of ZrO_2 . Small single crystals were grown and found to adopt a pseudo-orthorhombic unit cell with $a = 9.941$, $b = 11.482$, and $c = 23.528$ Å.^[15] A later preliminary structural refinement of these crystals^[16] suggested that the actual composition was $\text{Ba}_{11}(\text{Ti,Zr})_{28}\text{O}_{66}\text{O}$, with the 67th oxygen in a position not located.^[9] Subsequent chemical analyses of these crystals using an ion-microprobe^[12] found about 0.25 wt% Al_2O_3 in addition to the expected amounts of BaO , TiO_2 , and ZrO_2 , although no alumina had been intentionally added. A comparison of

the pseudo-orthorhombic unit cell for these crystals as well as the proposed crystal structure (Figure 10 in ref.^[9]) reveals that they are essentially identical to those for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ reported here (Table 1, Figure 3). Apparently, the phase “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ” is actually a ternary compound stabilized by a relatively small content of a trivalent metal such as Fe^{3+} or Al^{3+} . Note that the A:B:O A-cation:B-cation:oxygen ratios are very similar for all of these compositions. As previously noted,^[10] Jonker and Kwestroo^[4] “wet-milled” their specimens twice, once before and once after calcination, and this process is notorious for the introduction of impurities from the ceramic spheres used for milling. To confirm this, specimens with compositions $\text{Ba}_{11}\text{Al}_2\text{Ti}_{26}\text{O}_{66}$, $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Sn}_2\text{O}_{66}$, and $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Zr}_2\text{O}_{66}$ were synthesized (1250 °C, 1275 °C, and 1275 °C, respectively), and were all found to contain the $\text{A}_{11}\text{B}_{28}\text{O}_{66+x}$ -type phase as the major product (>95%) [the phase was not observed to form in compositions with Mg^{2+} or Zn^{2+} (i.e. $\text{Ba}_{11}\text{MgTi}_{27}\text{O}_{66}$ or $\text{Ba}_{11}\text{ZnTi}_{27}\text{O}_{66}$)].

Magnetic and Dielectric Properties of $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$

The inverse magnetic susceptibility vs. temperature measured for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ is shown in Figure 5. The behavior is paramagnetic, with highly dilute magnetic species, and follows the Curie–Weiss law above about 75 K. The slight curvature below 75 K and the extrapolated temperature intercept of –25 K indicate weak cooperative interactions between Fe spins that are, overall, antiferromagnetic in nature. Application of the Curie–Weiss formalism to the near-linear data above 75 K yields an effective moment of $5.2 \mu_B/\text{Fe}$, which is similar to the free-ion spin-only moment for high-spin Fe^{3+} ($S = 5/2$; $5.9 \mu_B/\text{Fe}$).

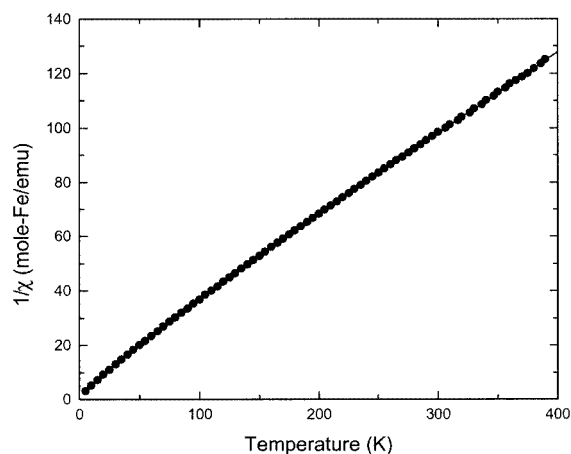


Figure 5. Inverse magnetic susceptibility vs. temperature for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$; the slight curvature below 75 K and the extrapolated temperature intercept of –25 K indicate weak cooperative interactions among the magnetic species that are overall antiferromagnetic in nature; the effective magnetic moment calculated from the essentially linear data above 75 K is consistent with the presence of high-spin Fe^{3+}

A corrected relative permittivity of 55, and a dielectric loss tangent, $\tan\delta$, of $7.7(\pm 0.3) \times 10^{-4}$ were obtained for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$. Both values were essentially independent of frequency between 5.0 and 5.5 GHz. Using Shannon's values for atomic polarizabilities and the Clausius–Mosotti theory,^[18] a dielectric constant of 45 was estimated for this compound, a value in fair agreement with the experimental results given the uncertainty of the porosity measurement.

Conclusions

The crystal structure of the new ternary oxide $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ has been characterized by single-crystal and powder X-ray diffraction methods. The compound adopts a monoclinic $10L$ structure built from close-packed (cp) $[\text{O}(\text{Ba},\text{O})]$ layers with a stacking sequence $(cchhc)_2$. Octahedral sites are occupied by Fe^{3+} and Ti^{4+} . The structure features vertex-, edge-, and face-sharing of the $[\text{Ti}(\text{Fe})\text{O}_6]$ octahedra. $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ is built from the same three types of layers as the neighboring $8L$ compound $\text{Ba}_4\text{Fe}_2\text{Ti}_{10}\text{O}_{27}$. As observed for the binary barium polytitanates, some of the cp layers also feature vacant sites between two Ba ions which substitute for three oxygens in a row. The single-crystal structural results for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ indicate that one of these sites is partially occupied, and that the basic structural formula for this phase is $\text{A}_{11}\text{B}_{28}\text{O}_{66+x}$, where $\text{A} = \text{Ba}^{2+}$ and $\text{B} = \text{Ti}^{4+}$ substituted with a small amount (<10 atom%) of a trivalent metal, which determines the value of x . The stoichiometry, or x -value, of the phase prepared in polycrystalline form at 1260 °C is slightly different from that of the crystals grown by cooling from 1390 °C, and occurs between the compositions $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ and $\text{Ba}_{11}\text{Fe}_2\text{Ti}_{26}\text{O}_{66}$. This variation in oxygen content and Fe/Ti ratio suggests that the partial occupancy site factor depends on the synthesis temperature. Interestingly, the phase apparently forms as a point compound in the subsolidus, with little or no solid solution formation. Analysis of the structural details using the bond valence method revealed one unusual B site, which was relatively expanded and isolated within the structure. This site is probably preferentially occupied by the larger, lower-valent Fe^{3+} cations. The crystal structure obtained here for $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ is essentially identical to that proposed earlier for “ Zr^{4+} -stabilized $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ”^[9,15] crystals. The phase originally reported as “ $\text{Ba}_2\text{Ti}_5\text{O}_{12}$ ”^[10] is actually a ternary compound, which forms upon addition (either deliberately or inadvertently via, e.g., ball-milling) of a trivalent ion such as Fe^{3+} or Al^{3+} . The specimens $\text{Ba}_{11}\text{Al}_2\text{Ti}_{26}\text{O}_{66}$, $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Sn}_2\text{O}_{66}$, and $\text{Ba}_{11}\text{Al}_2\text{Ti}_{24}\text{Zr}_2\text{O}_{66}$ were also prepared, and were found to form the $\text{A}_{11}\text{B}_{28}\text{O}_{66+x}$ -type phase.

$\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ was found to exhibit paramagnetic behavior above about 75 K. Application of the Curie–Weiss formalism to the $1/\chi$ vs. T data above 75 K yielded an effective moment consistent with the presence of high-spin Fe^{3+} , and a negative Weiss constant (about -25 K) indicating weak cooperative magnetic interactions that are overall antiferromagnetic. The relative permittivity and dielectric

loss tangent of a sintered polycrystalline disk were measured at 5.33 GHz, yielding values (corrected for theoretical density) of 55 and $7.7(\pm 0.3) \times 10^{-4}$, respectively.

Experimental Section

Dark reddish-brown polycrystalline $\text{Ba}_{11}\text{FeTi}_{27}\text{O}_{66.5}$ was prepared by solid-state reaction of reagent-grade BaCO_3 , Fe_2O_3 , and phosphate-free TiO_2 in air in the molar ratio $\text{BaO}:\text{Fe}_2\text{O}_3:\text{TiO}_2 = 0.28:0.02:0.70$. Prior to each heating, the sample was ground for 15–20 min with an agate pestle and mortar, pressed into pellets, and placed on sacrificial powder of the same composition in an alumina combustion boat. After an initial overnight calcine at 1000 °C, four one-week heatings at 1260 °C resulted in a well crystallized, nearly single-phase* sample (10-g scale) [* the pattern contains two weak peaks ($I_{\text{rel}} < 0.8\%$), which can be attributed to $\text{Ba}_6\text{Ti}_{17}\text{O}_{40}$. Detection of small amounts of the phase $\text{Ba}_4\text{Fe}_2\text{Ti}_{10}\text{O}_{27}$ was precluded by peak overlap with the title phase]. Dark red platelet-like single crystals were obtained by slowly cooling (1 °C/h) a slightly non-stoichiometric liquid ($\text{BaO}:\text{Fe}_2\text{O}_3:\text{TiO}_2 = 0.28:0.02:0.70$) from 1390 °C to 1260 °C, followed by step-cooling to 800 °C, and then air-quenching. Samples were contained in platinum capsules (2.6 mm i.d.) open to the air. Crystals were harvested from dark prismatic chunks obtained by breaking the crystallized reaction mass, and were then characterized by the precession camera method (Zr-filtered Mo-K_α radiation) to assess quality, approximate cell parameters, and space group.

Single-crystal X-ray diffraction data were obtained with an Enraf–Nonius CAD-4 diffractometer (certain commercial equipment is identified in order to adequately specify the experimental procedure; recommendation or endorsement by the National Institute of Standards and Technology is not therein implied) using graphite-monochromated Mo-K_α radiation. A Gaussian integration absorption correction was applied to all measured intensities. Calculations were carried out using the NRCVAX^[19] suite of programs. The structure was solved using Patterson and Fourier methods. Octahedral sites were assumed to be randomly occupied by Ti^{4+} and a small amount (< 10 atom%) of Fe^{3+} . Since the contrast between Fe and Ti by X-ray diffraction is small, no attempt was made to directly refine their relative occupancies, and the scattering factor for Ti was used for atoms occupying the octahedral sites. The overall stoichiometry of the phase reported here was then deduced by simultaneous consideration of the following: 1) the disappearing phase method,^[3] 2) charge balance in the presence of trivalent iron and tetravalent titanium, and 3) the degree of site occupancy of oxygen as indicated by the single-crystal X-ray diffraction data.

X-ray powder diffraction data were collected using a Philips diffractometer equipped with incident soller slits, theta-compensating slits, a 0.2 mm receiving slit, a graphite monochromator, and a scintillation detector. Data were collected at ambient temperature using Cu-K_α radiation with a 0.01° 2θ step size, and a 2 s count time. The sample was mounted in a wetted glass slide. Intensity data, measured as relative peak heights above the background of hand-picked peaks, were obtained using the Jade 6 program suite. The observed 20 line positions reported here were corrected using SRM 660, LaB_6 ,^[20] as an external calibrating agent. The unit cell was refined from the corrected powder diffraction data (20 values; $\text{Cu-K}_\alpha = 1.540593$ Å) using the least-squares program CELLSVD.^[21] During the refinement, the indices of selected peaks were fixed in

accordance with the powder pattern calculated using the results of the single-crystal X-ray study and the Jade program jPOWD.

The conditions used for the single-crystal structure determination are given in Table 4. Further details of the crystal-structure investigation may be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany, on quoting depository number CSD-413889.

Table 4. Experimental details for the single-crystal X-ray diffraction study of Ba₁₁FeTi₂₇O_{66.5}

Crystal system	Monoclinic
Space group	<i>C</i> 2/m
Cell parameters	<i>a</i> = 23.321(2) Å <i>b</i> = 11.384(1) Å <i>c</i> = 9.847(1) Å β = 90.09(1) ° <i>V</i> = 2614.2(8) Å ³
<i>Z</i>	2
$\rho_{\text{calcd.}}$ [g cm ⁻³]	4.98
Crystal dimensions [mm]	0.1 × 0.1 × 0.03
μ [mm ⁻¹]	12.22
2 θ (max)	60°
Scan mode	T-scan
Reflections measured	15759
Unique reflections	3954
Reflections observed	1656
Transmission factors	0.15–0.31
Number of parameters	157
<i>R</i> _F	0.088
<i>wR</i> _F	0.065

Magnetization was measured using a Quantum Design SQUID (Superconducting Quantum Interference Device) magnetometer. Measurements were made on a polycrystalline sample in a field of 0.5 Tesla between 400 K and 5 K. Magnetization measurements were carried out in fields up to 4 Tesla.

Permittivity and dielectric loss tangent measurements were performed at 5.33 GHz under ambient conditions using a sintered polycrystalline disk, approximately 11 mm in diameter and 87% dense. As described previously,^[6,7] a dielectric resonator technique utilizing higher-order TE_{0_{xy}} modes was used that allows the measurement of dielectric properties of a single sample at several contiguous frequency sub-bands. Corrections to the crystallographic density were incorporated by estimating the pore volume of the sample (from size and mass), and applying the Bruggeman effective medium formulation^[22] for a two-phase composite. Uncertainties in the measurement of pore volume suggest a realistic accuracy limit for the permittivity values of ±10%, confirmed by measure-

ments of similarly dense test samples such as LaAlO₃,^[6] for which reliable single-crystal data are available.

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